

## Global Solvability and Asymptotic Behaviour of a Hyperbolic Problem With Acoustic Boundary Conditions

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### 1. Introduction

This article presents a study of global solvability and asymptotic behaviour of solutions to the mixed problem for the nonlinear hyperbolic equation with nonlinear damping and acoustic boundary conditions:

$$(1.1) \quad u_{tt} - a(u)u_{xx} + g(u_t) = 0, \quad \text{in } Q = (0, 1) \times (0, T),$$

$$(1.2) \quad u(0, t) = 0, \quad t \in (0, T),$$

$$(1.3) \quad -\rho u_t(1, t) = \delta_{tt}(t) + c_1 \delta_t(t) + c_0 \delta(t), \quad t \in (0, T),$$

$$(1.4) \quad u_x(1, t) = \delta_t(t), \quad t \in (0, T),$$

$$(1.5) \quad u(x, 0) = u_0(x), \quad u_t(x, 0) = u_1(x), \quad x \in (0, 1), \quad \delta(0) = \delta_0;$$

where  $\rho$ ,  $T$ ,  $c_0$ ,  $c_1$  are positive constants,  $a : \mathbb{R} \rightarrow \mathbb{R}$  and  $g : \mathbb{R} \rightarrow \mathbb{R}$  are differentiable functions.

In [3], J. T. Beale and S. I. Rosencrans considered acoustic boundary conditions for a gas, undergoing small irrotational perturbations from the rest, in a domain  $\Omega \subset \mathbb{R}^3$  with a smooth compact boundary  $\partial\Omega = \Gamma$ . They assumed that the boundary  $\Gamma$  is nonporous and locally reacting in the sense that wave motion along the boundary is negligible. Therefore, if  $u(x, t)$  is the velocity potential and  $\delta(x, t)$  denotes the normal displacement of a point  $x \in \Gamma$  at time  $t$ , they must satisfy the equations:

$$(1.6) \quad u_{tt} - \Delta u = 0, \quad \text{in } \Omega \times (0, T);$$

$$(1.7) \quad \delta_{tt} - c_1 \delta_t + c_0 \delta = -\rho u_t, \quad \text{on } \Gamma \times (0, T);$$

$$(1.8) \quad \delta_t = \frac{\partial u}{\partial \nu}, \quad \text{on } \Gamma \times (0, T);$$

here  $\nu$  is the unit outer normal,  $c_1 = \frac{d}{m}$ ,  $c_0 = \frac{k}{m}$ ,  $\rho = \frac{\rho_0}{m}$ ; where  $m$  is the mass per unit area on the boundary,  $d$  is the resistivity of the boundary,  $k$  is the spring constant and  $\rho_0$  is the unperturbed density of the gas.

The mixed problem for (1.6) – (1.8) was considered by J. T. Beale which, using semigroup techniques, carried out a detailed analysis in bounded domain [1] and exterior domains [2]. Beale's idea was to work with an equivalent initial value problem in the form  $u_t = Au$  where  $A$  is an operator on a suitable Hilbert space  $H$ . For bounded domains, in [1], Beale gave a description of the spectrum of the operator  $A$  and proved that there is no uniform rate of decay for solutions to the mixed problem for (1.6) – (1.8).

Recently, C. L. Frota and J. A. Goldstein [4] studied acoustic boundary conditions (1.7), (1.8) for the nonlinear Carrier equation with a nonlinear damping

$$(1.9) \quad u_{tt} - M\left(\int_{\Omega} u^2 dx\right)\Delta u + |u_t|^\alpha u_t = 0.$$

The authors proved results on global solvability, uniqueness, regularity and continuous dependence of solutions on the parameters when the boundary  $\Gamma$  is made up of two disjoint pieces  $\Gamma_0, \Gamma_1$ , each having nonempty interior. Acoustic boundary conditions (1.7), (1.8) were imposed for  $(x, t) \in \Gamma_1 \times (0, T)$  and the Dirichlet boundary condition  $u(x, t) = 0$  was imposed for  $(x, t) \in \Gamma_0 \times [0, T]$ . Moreover, for dimensions  $n = 2$  or  $n = 3$ , global solvability and uniqueness for the mixed problem for (1.9) with acoustic boundary conditions on the whole boundary  $\Gamma$  were established. This was obtained as a limit case when the positive measure of  $\Gamma_0$  shrink to zero.

Global solvability and stability of the energy of the mixed problem for (1.9) with the Dirichlet boundary conditions was proved in [5].

In the present paper we consider the nonlinear wave equation (1.1) with local nonlinearities. The function  $a(u)$  depends on a solution while the function  $M(\int_{\Omega} u^2 dx)$  in (1.9) depends on the  $L^2$ -norm of it. This difference makes the study of (1.1) – (1.5) more complicated and we restrict it only to the one-dimensional case. Exponential and algebraic energy decay are proved under some appropriate assumptions on the damping  $g(u_t)$  and when  $c_0$  is sufficiently

large.

## 2. Existence and uniqueness

In this section we prove the existence and uniqueness of global solutions to (1.1) – (1.5). We denote the inner product and norm in  $L^2(0, 1)$  by

$$(u, v) = \int_0^1 u(x)v(x) dx, \quad \|u\| = \left( \int_0^1 |u(x)|^2 dx \right)^{\frac{1}{2}},$$

and we define  $V = \{z \in H^1(0, 1); z(0) = 0\}$  and the norm in  $V$  by

$$\|z\|_V = \|z_x\| \quad z \in V.$$

The functional spaces are standard, see in Lions and Magenes [5].

**Theorem 2.1.** *Suppose that the functions  $a$  and  $g$  satisfy*

$$(2.1) \quad 0 < a_0 \leq a(\lambda),$$

$$(2.2) \quad |a'(\lambda)| \leq k_0 a(\lambda),$$

$$(2.3) \quad |g(\xi)| \leq k_1 (|\xi| + |\xi|^{\alpha+1}),$$

$$(2.4) \quad k_2 |\xi|^{\alpha+2} \leq g(\xi) \xi,$$

$$(2.5) \quad k_3 |\xi|^\alpha \leq g'(\xi),$$

for all  $\lambda, \xi \in \mathbb{R}$ , where  $k_i; i = 0, 1, 2, 3$ ; are positive constants and  $\alpha > 1$ . Let  $c_0, \rho$  be positive constants and  $c_1 \in \mathbb{R}$ . Given  $\delta_0 \in \mathbb{R}$ ,  $u_0 \in V \cap H^2(0, 1)$  and  $u_1 \in V \cap L^{2\alpha+2}(0, 1)$ , there exists a unique pair of functions  $(u(x, t), \delta(t))$  which is a solution to (1.1) – (1.5) from the class:

$$u, u' \in L^\infty(0, T; C([0, 1])), \quad u'' \in L^\infty(0, T; L^2(0, 1));$$

$$\delta \in C^1([0, T]), \quad \delta'' \in L^\infty(0, 1).$$

**Proof.** Let  $\{w_j\}$  be a basis of  $V$ , orthonormal in  $L^2(0, 1)$ . For each  $\epsilon \in (0, 1]$  and  $m \in \mathbb{N}$  we consider  $\delta_{\epsilon m} : [0, t] \rightarrow \mathbb{R}$  and  $u_{\epsilon m}(x, t) = \sum_{j=1}^m \varphi_{\epsilon m j}(t) w_j(x)$ ,  $x \in$

$(0, 1)$ ,  $t \in [0, T_{\epsilon m}]$ , which are solutions to the following approximate problem

$$\left\{ \begin{array}{l} \left( \left( \epsilon + \frac{1}{a(u_{\epsilon m}(t))} \right) u_{\epsilon m}''(t), w_j \right) + \left( \frac{\partial u_{\epsilon m}}{\partial x}(t), \frac{\partial w_j}{\partial x} \right) - \delta'_{\epsilon m}(t) w_j(1) \\ + \left( \frac{g(u'_{\epsilon m}(t))}{a(u_{\epsilon m}(t))}, w_j \right) = 0, \quad 1 \leq j \leq m, \\ \rho u'_{\epsilon m}(1, t) + \delta''_{\epsilon m}(t) + c_1 \delta'_{\epsilon m}(t) + c_0 \delta_{\epsilon m}(t) = 0, \\ u_{\epsilon m}(0) = u_{0m}, \quad u'_{\epsilon m}(0) = u_{1m}, \quad \delta_{\epsilon m}(0) = \delta_0, \quad \delta'_{\epsilon m}(0) = \frac{\partial u_{0m}}{\partial x}(1); \end{array} \right.$$

where  $\delta_0 \in \mathbb{R}$ ,  $u_{0m} = \sum_{j=1}^m (u_0, w_j) w_j$  and  $u_{1m} = \sum_{j=1}^m (u_1, w_j) w_j$ .

Since  $\epsilon > 0$ , this approximate problem can be transformed into a normal system of ODE for the unknown functions  $\delta_{\epsilon m}$  and  $\varphi_{\epsilon m j}$ ,  $1 \leq j \leq m$ . Therefore, it has local solutions, defined at some interval  $[0, T_{\epsilon m}]$ , which satisfy the following approximate equations:

$$(2.6) \quad \left( \left( \epsilon + \frac{1}{a(u_{\epsilon m}(t))} \right) u_{\epsilon m}''(t), v \right) + \left( \frac{\partial u_{\epsilon m}}{\partial x}(t), \frac{\partial v}{\partial x} \right) - \delta'_{\epsilon m}(t) v(1) \\ + \left( \frac{g(u'_{\epsilon m}(t))}{a(u_{\epsilon m}(t))}, v \right) = 0, \quad \text{for all } v \in V_m = [w_1, \dots, w_m],$$

$$(2.7) \quad -\rho u'_{\epsilon m}(1, t) = \delta''_{\epsilon m}(t) + c_1 \delta'_{\epsilon m}(t) + c_0 \delta_{\epsilon m}(t).$$

**Estimate 1.** Taking  $v = 2u'_{\epsilon m}(t)$  in (2.6) and multiplying (2.7) by  $2\delta'_{\epsilon m}(t)$ , we have

$$\begin{aligned} & \frac{d}{dt} \left[ \int_0^1 \left( \epsilon + \frac{1}{a(u_{\epsilon m}(x, t))} \right) (u'_{\epsilon m}(x, t))^2 dx + \|u_{\epsilon m}(t)\|_V^2 \right] - 2\delta'_{\epsilon m}(t) u'_{\epsilon m}(1, t) \\ & + 2 \int_0^1 \frac{g(u'_{\epsilon m}(x, t)) u'_{\epsilon m}(x, t)}{a(u_{\epsilon m}(x, t))} dx = - \int_0^1 \frac{a'(u_{\epsilon m}(x, t)) (u'_{\epsilon m}(x, t))^3}{(a(u_{\epsilon m}(x, t)))^2} dx; \\ & -2\delta'_{\epsilon m}(t) u'_{\epsilon m}(1, t) = \frac{d}{dt} \left[ \frac{1}{\rho} (\delta'_{\epsilon m}(t))^2 + \frac{c_0}{\rho} (\delta_{\epsilon m}(t))^2 \right] + \frac{2c_1}{\rho} (\delta'_{\epsilon m}(t))^2. \end{aligned}$$

From here we get

$$\begin{aligned} & \frac{d}{dt} \left[ \int_0^1 \left( \epsilon + \frac{1}{a(u_{\epsilon m}(x, t))} \right) (u'_{\epsilon m}(x, t))^2 dx + \|u_{\epsilon m}(t)\|_V^2 + \frac{1}{\rho} (\delta'_{\epsilon m}(t))^2 \right. \\ & \left. + \frac{c_0}{\rho} (\delta_{\epsilon m}(t))^2 \right] + 2 \int_0^1 \frac{g(u'_{\epsilon m}(x, t)) u'_{\epsilon m}(x, t)}{a(u_{\epsilon m}(x, t))} dx \\ & = - \int_0^1 \frac{a'(u_{\epsilon m}(x, t)) (u'_{\epsilon m}(x, t))^3}{(a(u_{\epsilon m}(x, t)))^2} dx - \frac{2c_1}{\rho} (\delta'_{\epsilon m}(t))^2. \end{aligned}$$

This equality, (2.2) and (2.4) give

$$\begin{aligned} & \frac{d}{dt} \left[ \int_0^1 \left( \epsilon + \frac{1}{a(u_{\epsilon m}(x, t))} \right) (u'_{\epsilon m}(x, t))^2 dx + \|u_{\epsilon m}(t)\|_V^2 + \frac{1}{\rho} (\delta'_{\epsilon m}(t))^2 \right. \\ & \left. + \frac{c_0}{\rho} (\delta_{\epsilon m}(t))^2 \right] + 2k_2 \int_0^1 \frac{|u'_{\epsilon m}(x, t)|^{\alpha+2}}{a(u_{\epsilon m}(x, t))} dx \leq k_0 \int_0^1 \frac{|u'_{\epsilon m}(x, t)|^3}{a(u_{\epsilon m}(x, t))} dx \\ & + \frac{2|c_1|}{\rho} (\delta'_{\epsilon m}(t))^2. \end{aligned}$$

Using Young's inequality for all  $\eta > 0$ , we obtain

$$\int_0^1 \frac{|u'_{\epsilon m}(x, t)|^3}{a(u_{\epsilon m}(x, t))} dx \leq \eta \int_0^1 \frac{|u'_{\epsilon m}(x, t)|^{\alpha+2}}{a(u_{\epsilon m}(x, t))} dx + c(\eta) \int_0^1 \frac{|u'_{\epsilon m}(x, t)|^2}{a(u_{\epsilon m}(x, t))} dx.$$

Choosing  $\eta$  sufficiently small, we find

$$\begin{aligned} & \frac{d}{dt} \left[ \int_0^1 \left( \epsilon + \frac{1}{a(u_{\epsilon m}(x, t))} \right) (u'_{\epsilon m}(x, t))^2 dx + \|u_{\epsilon m}(t)\|_V^2 + \frac{1}{\rho} (\delta'_{\epsilon m}(t))^2 \right. \\ & \left. + \frac{c_0}{\rho} (\delta_{\epsilon m}(t))^2 \right] + A_1 \int_0^1 \frac{|u'_{\epsilon m}(x, t)|^{\alpha+2}}{a(u_{\epsilon m}(x, t))} dx \leq A_2 \int_0^1 \frac{|u'_{\epsilon m}(x, t)|^2}{a(u_{\epsilon m}(x, t))} dx \\ & + A_3 (\delta'_{\epsilon m}(t))^2; \end{aligned}$$

here and in expressions to follow,  $A_i$ ,  $i \in \mathbb{N}$ , denotes a positive constant which does not depend on  $\epsilon$ ,  $m$  and  $t$ . Integrating this inequality from 0 to  $t \leq T_{\epsilon m}$ , we obtain

$$\begin{aligned} & \int_0^1 \left( \epsilon + \frac{1}{a(u_{\epsilon m}(x, t))} \right) (u'_{\epsilon m}(x, t))^2 dx + \|u_{\epsilon m}(t)\|_V^2 + \frac{1}{\rho} (\delta'_{\epsilon m}(t))^2 + \frac{c_0}{\rho} (\delta_{\epsilon m}(t))^2 \\ & + A_1 \int_0^t \int_0^1 \frac{|u'_{\epsilon m}(x, \tau)|^{\alpha+2}}{a(u_{\epsilon m}(x, \tau))} dx d\tau \leq A_4 + A_2 \int_0^t \int_0^1 \frac{|u'_{\epsilon m}(x, \tau)|^2}{a(u_{\epsilon m}(x, \tau))} dx d\tau \\ & + A_3 \int_0^t (\delta'_{\epsilon m}(\tau))^2 d\tau. \end{aligned}$$

Gronwall's inequality yields

$$\begin{aligned} (2.8) \quad & \int_0^1 \left( \epsilon + \frac{1}{a(u_{\epsilon m}(x, t))} \right) (u'_{\epsilon m}(x, t))^2 dx + \|u_{\epsilon m}(t)\|_V^2 + \frac{1}{\rho} (\delta'_{\epsilon m}(t))^2 \\ & + \frac{c_0}{\rho} (\delta_{\epsilon m}(t))^2 + A_1 \int_0^t \int_0^1 \frac{|u'_{\epsilon m}(x, \tau)|^{\alpha+2}}{a(u_{\epsilon m}(x, \tau))} dx d\tau \leq A_5. \end{aligned}$$

Taking into account that  $V \hookrightarrow H^1(0, 1) \hookrightarrow C([0, 1])$  and  $a$  is a continuous function, we can see that there exists a constant  $M$ , independent of  $\epsilon$ ,  $m$  and  $t$ , such that

$$(2.9) \quad 0 < a_0 \leq a(u_{\epsilon m}(x, t)) \leq M; \quad \text{for all } x \in (0, 1), \quad t \in [0, T_{\epsilon m}].$$

From (2.8) and (2.9), we have

$$(2.10) \quad \|u'_{\epsilon m}(t)\|^2 + \|u_{\epsilon m}(t)\|^2 + (\delta'_{\epsilon m}(t))^2 + (\delta_{\epsilon m}(t))^2 + \int_0^t \|u'_{\epsilon m}(\tau)\|_{\alpha+2}^{\alpha+2} d\tau \leq A_6.$$

This is the first estimate.

**Estimate 2.** First of all we put  $t = 0$  in (2.6) and (2.7). It follows

$$\begin{aligned} \left( \left( \epsilon + \frac{1}{a(u_{0m})} \right) u''_{\epsilon m}(0), v \right) - \left( \frac{\partial^2 u_{0m}}{\partial x^2}, v \right) + \left( \frac{g(u_{1m})}{a(u_{0m})}, v \right) &= 0, \\ \delta''_{\epsilon m}(0) &= -\rho u_{1m}(1) - c_1 \frac{\partial u_{0m}}{\partial x}(1) - c_0 \delta_0. \end{aligned}$$

Taking  $v = u''_{\epsilon m}(0)$  in this equality, we find constants  $A_7$  and  $A_8$ , such that

$$(2.11) \quad \|u''_{\epsilon m}(0)\| \leq A_7 \quad \text{and} \quad |\delta''_{\epsilon m}(0)| \leq A_8.$$

Differentiating (2.6) and (2.7) with respect to  $t$ , we have

$$(2.12) \quad \begin{aligned} \left( \left( \epsilon + \frac{1}{a(u_{\epsilon m}(t))} \right) u'''_{\epsilon m}(t), v \right) + \left( \frac{\partial u'_{\epsilon m}}{\partial x}(t), \frac{\partial v}{\partial x} \right) - \delta''_{\epsilon m}(t)v(1) \\ \left( \frac{g'(u'_{\epsilon m}(t))u''_{\epsilon m}(t)}{a(u_{\epsilon m}(t))}, v \right) - \left( \frac{g(u'_{\epsilon m}(t))u'_{\epsilon m}(t)a'(u_{\epsilon m}(t))}{[a(u_{\epsilon m}(t))]^2}, v \right) \\ - \left( \frac{a'(u_{\epsilon m}(t))u'_{\epsilon m}(t)u''_{\epsilon m}(t)}{[a(u_{\epsilon m}(t))]^2}, v \right) &= 0, \end{aligned}$$

$$(2.13) \quad \delta'''_{\epsilon m}(t) + c_1 \delta''_{\epsilon m}(t) + c_0 \delta'_{\epsilon m}(t) = -\rho u''_{\epsilon m}(1, t).$$

Taking  $v = 2u''_{\epsilon m}(t)$  in (2.12) and multiplying (2.13) by  $2\delta''_{\epsilon m}(t)$ , we obtain

$$\begin{aligned} \frac{d}{dt} \left[ \int_0^1 \left( \epsilon + \frac{1}{a(u_{\epsilon m}(x, t))} \right) |u''_{\epsilon m}(x, t)|^2 dx + \|u'_{\epsilon m}(t)\|_V^2 + \frac{1}{\rho} (\delta''_{\epsilon m}(t))^2 \right. \\ \left. + \frac{c_0}{\rho} (\delta'_{\epsilon m}(t))^2 \right] + 2 \int_0^1 \frac{g'(u'_{\epsilon m}(x, t)) |u''_{\epsilon m}(x, t)|^2}{a(u_{\epsilon m}(x, t))} dx = -\frac{2c_1}{\rho} (\delta''_{\epsilon m}(t))^2 \\ + \int_0^1 \frac{a'(u_{\epsilon m}(x, t)) u'_{\epsilon m}(x, t) |u''_{\epsilon m}(x, t)|^2}{[a(u_{\epsilon m}(x, t))]^2} dx \\ + 2 \int_0^1 \frac{a'(u_{\epsilon m}(x, t)) g(u'_{\epsilon m}(x, t)) u'_{\epsilon m}(x, t) u''_{\epsilon m}(x, t)}{[a(u_{\epsilon m}(x, t))]^2} dx. \end{aligned}$$

For all  $\eta > 0$ , Young's inequality gives

$$\begin{aligned} I_1 &= \int_0^1 \frac{a'(u_{\epsilon m}(x, t)) u'_{\epsilon m}(x, t) |u''_{\epsilon m}(x, t)|^2}{[a(u_{\epsilon m}(x, t))]^2} dx \\ &\leq \eta \int_0^1 \frac{|u'_{\epsilon m}(x, t)|^\alpha |u''_{\epsilon m}(x, t)|^2}{a(u_{\epsilon m}(x, t))} dx + C(\eta) \int_0^1 \frac{|u''_{\epsilon m}(x, t)|^2}{a(u_{\epsilon m}(x, t))} dx; \end{aligned}$$

$$\begin{aligned} I_2 &= 2 \int_0^1 \frac{a'(u_{\epsilon m}(x, t)) g(u'_{\epsilon m}(x, t)) u'_{\epsilon m}(x, t) u''_{\epsilon m}(x, t)}{[a(u_{\epsilon m}(x, t))]^2} dx \\ &\leq 2k_0 k_1 \left( \int_0^1 \frac{|u'_{\epsilon m}(x, t)|^2 |u''_{\epsilon m}(x, t)|}{a(u_{\epsilon m}(x, t))} dx + \int_0^1 \frac{|u'_{\epsilon m}(x, t)|^{\alpha+2} |u''_{\epsilon m}(x, t)|}{a(u_{\epsilon m}(x, t))} dx \right) \\ &\leq \eta \int_0^1 \frac{|u'_{\epsilon m}(x, t)|^\alpha |u''_{\epsilon m}(x, t)|^2}{a(u_{\epsilon m}(x, t))} dx + C(\eta) \|u'_{\epsilon m}(t)\|_V^2 \int_0^1 \frac{|u'_{\epsilon m}(x, t)|^{\alpha+2}}{a(u_{\epsilon m}(x, t))} dx \\ &\quad + 2k_0 k_1 \int_0^1 \frac{|u''_{\epsilon m}(x, t)|^2}{a(u_{\epsilon m}(x, t))} dx + A_{11}; \end{aligned}$$

$$I_3 = 2 \int_0^1 \frac{g'(u'_{\epsilon m}(x, t)) |u''_{\epsilon m}(x, t)|^2}{a(u_{\epsilon m}(x, t))} dx \geq 2k_3 \int_0^1 \frac{|u'_{\epsilon m}(x, t)|^\alpha |u''_{\epsilon m}(x, t)|^2}{a(u_{\epsilon m}(x, t))} dx;$$

here we have used (2.3) and (2.4).

Therefore,

$$\begin{aligned} &\frac{d}{dt} \left[ \int_0^1 \left( \epsilon + \frac{1}{a(u_{\epsilon m}(x, t))} \right) |u''_{\epsilon m}(x, t)|^2 dx + \|u'_{\epsilon m}(t)\|_V^2 + \frac{1}{\rho} (\delta''_{\epsilon m}(t))^2 \right. \\ &\quad \left. + \frac{c_0}{\rho} (\delta'_{\epsilon m}(t))^2 \right] + 2k_3 \int_0^1 \frac{|u'_{\epsilon m}(x, t)|^\alpha |u''_{\epsilon m}(x, t)|^2}{a(u_{\epsilon m}(x, t))} dx \leq \frac{2|c_1|}{\rho} (\delta''_{\epsilon m}(t))^2 \\ &\quad + \eta \int_0^1 \frac{|u'_{\epsilon m}(x, t)|^\alpha |u''_{\epsilon m}(x, t)|^2}{a(u_{\epsilon m}(x, t))} dx + C(\eta) \|u'_{\epsilon m}(t)\|_V^2 \int_0^1 \frac{|u'_{\epsilon m}(x, t)|^{\alpha+2}}{a(u_{\epsilon m}(x, t))} dx \\ &\quad + \int_0^1 \frac{|u''_{\epsilon m}(x, t)|^2}{a(u_{\epsilon m}(x, t))} dx + A_{11}. \end{aligned}$$

Choosing  $\eta > 0$  sufficiently small we have

$$\begin{aligned} &\frac{d}{dt} \left[ \int_0^1 \left( \epsilon + \frac{1}{a(u_{\epsilon m}(x, t))} \right) |u''_{\epsilon m}(x, t)|^2 dx + \|u'_{\epsilon m}(t)\|_V^2 + \frac{1}{\rho} (\delta''_{\epsilon m}(t))^2 \right. \\ &\quad \left. + \frac{c_0}{\rho} (\delta'_{\epsilon m}(t))^2 \right] + A_{12} \int_0^1 \frac{|u'_{\epsilon m}(x, t)|^\alpha |u''_{\epsilon m}(x, t)|^2}{a(u_{\epsilon m}(x, t))} dx \leq A_{13} \left( (\delta''_{\epsilon m}(t))^2 \right. \\ &\quad \left. + \int_0^1 \frac{|u''_{\epsilon m}(x, t)|^2}{a(u_{\epsilon m}(x, t))} dx + \|u'_{\epsilon m}(t)\|_V^2 \int_0^1 \frac{|u'_{\epsilon m}(x, t)|^{\alpha+2}}{a(u_{\epsilon m}(x, t))} dx + 1 \right). \end{aligned}$$

Integrating this, we get

$$\begin{aligned} \Psi_{\epsilon m}(t) &+ \int_0^t \int_0^1 \frac{|u'_{\epsilon m}(x, \tau)|^\alpha |u''_{\epsilon m}(x, \tau)|^2}{a(u_{\epsilon m}(x, \tau))} dx d\tau \\ &\leq A_{14} \left(1 + \int_0^t h_{\epsilon m}(\tau) \Psi_{\epsilon m}(\tau) d\tau\right), \end{aligned}$$

where

$$\begin{aligned} \Psi_{\epsilon m}(t) &= \int_0^1 \left(\epsilon + \frac{1}{a(u_{\epsilon m}(x, t))}\right) |u''_{\epsilon m}(x, t)|^2 dx + \|u'_{\epsilon m}(t)\|_V^2 \\ &+ (\delta''_{\epsilon m}(t))^2 + (\delta'_{\epsilon m}(t))^2, \end{aligned}$$

and

$$h_{\epsilon m}(t) = 1 + \int_0^1 \frac{|u'_{\epsilon m}(x, t)|^{\alpha+2}}{a(u_{\epsilon m}(x, t))} dx.$$

From (2.10) we can see that there exists a constant  $A_{15}$  independent on  $\epsilon, m, t$  such that

$$\int_0^t |h_{\epsilon m}(\tau)| d\tau \leq A_{15},$$

then  $h_{\epsilon m} \in L^1(0, T)$ . This and Gronwall's inequality give

$$(2.14) \quad \Psi_{\epsilon m}(t) \leq A_{16} e^{A_{14} \int_0^T h_{\epsilon m}(\tau) d\tau} = A_{17}, \quad \text{for all } t \in [0, T].$$

Hence

$$(2.15) \quad \|u''_{\epsilon m}(t)\|^2 + \|u'_{\epsilon m}(t)\|_V^2 + (\delta''_{\epsilon m}(t))^2 + (\delta'_{\epsilon m}(t))^2 \leq A_{18}, \quad \text{for all } t \in [0, T],$$

and the second estimate is proved.

Using estimates (2.10) and (2.15) we can pass to the limit as  $m \rightarrow \infty$  and  $\epsilon \rightarrow 0$  in the approximate problem to get

$$(2.16) \quad \left(\frac{u''(t)}{a(u(t))}, v\right) + \left(\frac{\partial u}{\partial x}(t), \frac{\partial v}{\partial x}\right) - \delta'(t)v(1) + \left(\frac{g(u'(t))}{a(u(t))}, v\right) = 0,$$

for all  $v \in V$ , a.e. on  $[0, T]$ .

$$(2.17) \quad \rho u'(1, t) + \delta''(t) + c_1 \delta'(t) + c_0 \delta(t) = 0, \quad \text{a.e. on } [0, T].$$

$$(2.18) \quad u(x, 0) = u_0(x), \quad u'(x, 0) = u_1(x), \quad \text{a.e. on } [0, 1].$$



$$(2.19) \quad \delta(0) = \delta_0, \quad \delta'(0) = \frac{\partial u_0}{\partial x}(1).$$

Taking in (2.16)  $v = \varphi \in \mathcal{D}(0, 1)$ , we have

$$\left( \frac{u''(t)}{a(u(t))} - \frac{\partial^2 u}{\partial x^2}(t) + \frac{g(u'(t))}{a(u(t))}, \varphi \right) = 0$$

which implies

$$u'' - a(u)u_{xx} + g(u') = 0 \quad \text{a.e. in } (0, 1) \times (0, T).$$

Uniqueness may be proved in the usual way. Theorem 2.1 is proved.

**Remark 2.1.** The result of Theorem 2.1 is valid even for  $\alpha = 1$  if we assume that  $k_2$  and  $k_3$  are sufficiently large.

### 3. Asymptotic behaviour

In this section, we shall prove the exponential and algebraic decay of the energy of (1.1) – (1.5). Therefore, instead of (2.3) and (2.4) we make the following stronger assumptions

$$(3.1) \quad |g(\xi)| \leq \bar{k}_1 (|\xi|^{\beta+1} + |\xi|^{\alpha+1}),$$

$$(3.2) \quad \bar{k}_2 (|\xi|^{\beta+2} + |\xi|^{\alpha+2}) \leq g(\xi)\xi;$$

for all  $\xi \in \mathbb{R}$ , with  $0 \leq \beta \leq 1$ , where

$$(3.3) \quad \bar{k}_2 > \left[ \frac{k_0(\alpha-1)}{2(\alpha-\beta)} \right]^{\frac{(\alpha-1)}{(\alpha-\beta)}} \left[ \frac{k_0(1-\beta)}{(\alpha-\beta)} \right]^{\frac{(1-\beta)}{(\alpha-\beta)}}.$$

**Theorem 3.1.** *Let  $(u, \delta)$  be a solution to the problem (1.1) – (1.5). Assume that  $c_1 > 0$  and the coefficient  $c_0$  in (1.3) satisfies the inequality*

$$(3.4) \quad \frac{\rho}{c_0} < \frac{1}{\sqrt{8\bar{c}}},$$

where  $\bar{c} > 0$  is the constant of the embedding  $\|\cdot\|_{C([0,1])} \leq \bar{c}\|\cdot\|_V$ . Then there are positive constants  $K$  and  $C$  such that

$$(3.5) \quad \left\{ \begin{array}{ll} h(t) \leq Ke^{-Ct}, \text{ for all } t \geq 0, & \text{if } \beta = 0 \\ \text{or} & \\ h(t) \leq C(1+t)^{\frac{-2}{\beta}}, \text{ for all } t \geq 0, & \text{if } 0 < \beta \leq 1 \end{array} \right.$$

here  $h(t) = \|u'(t)\|^2 + \|u(t)\|_V^2 + |\delta'(t)|^2 + |\delta(t)|^2$ .

**Proof.** Let us define a function  $\Phi : [0, \infty] \rightarrow \mathbb{R}$  by

$$(3.6) \quad \Phi(t) = \int_0^1 \frac{(u'(x, t))^2}{a(u(x, t))} dx + \|u(t)\|_V^2 + \frac{1}{\rho}(\delta'(t))^2 + \frac{c_0}{\rho}(\delta(t))^2.$$

Multiplying (1.1) by  $2u'(x, t)$  and integrating over  $(0, 1)$ , we obtain

$$\begin{aligned} & \frac{d}{dt} \left[ \int_0^1 \frac{(u'(x, t))^2}{a(u(x, t))} dx + \|u(t)\|_V^2 \right] - 2\delta'(t)u'(1, t) \\ & + 2 \int_0^1 \frac{g(u'(x, t))u'(x, t)}{a(u(x, t))} dx = - \int_0^1 \frac{a'(u(x, t))(u'(x, t))^3}{[a(u(x, t))]^2} dx. \end{aligned}$$

This and (1.3) give

$$(3.7) \quad \begin{aligned} \Phi'(t) + 2 \int_0^1 \frac{g(u'(x, t))u'(x, t)}{a(u(x, t))} dx + \frac{2c_1}{\rho}(\delta'(t))^2 \\ = - \int_0^1 \frac{a'(u(x, t))(u'(x, t))^3}{[a(u(x, t))]^2} dx. \end{aligned}$$

From (2.2) and (3.2), we find

$$\begin{aligned} \Phi'(t) + 2\bar{k}_2 \int_0^1 \frac{|u'(x, t)|^{\beta+2}}{a(u(x, t))} dx + 2\bar{k}_2 \int_0^1 \frac{|u'(x, t)|^{\alpha+2}}{a(u(x, t))} dx + \frac{2c_1}{\rho}(\delta'(t))^2 \\ \leq k_0 \int_0^1 \frac{|u'(x, t)|^3}{a(u(x, t))} dx. \end{aligned}$$

By Young's inequality

$$\begin{aligned} \Phi'(t) + \frac{2c_1}{\rho}(\delta'(t))^2 + (2\bar{k}_2 - \epsilon) \int_0^1 \frac{|u'(x, t)|^{\alpha+2}}{a(u(x, t))} dx \\ + \left[ 2\bar{k}_2 - \frac{k_0^{\frac{(\alpha-\beta)}{(\alpha-1)}}(\alpha-1)}{(\alpha-\beta)} \left( \frac{(\alpha-\beta)}{(1-\beta)} \epsilon \right)^{-\frac{(1-\beta)}{(\alpha-1)}} \right] \int_0^1 \frac{|u'(x, t)|^{\beta+2}}{a(u(x, t))} dx \leq 0, \end{aligned}$$

for all  $\epsilon > 0$ .

Choosing  $\epsilon = \bar{k}_2$ , and taking into account (3.3), we obtain

$$(3.8) \quad \Phi'(t) + \frac{2c_1}{\rho}(\delta'(t))^2 + A_1 \int_0^1 \frac{|u'(x, t)|^{\beta+2}}{a(u(x, t))} dx + \bar{k}_2 \int_0^1 \frac{|u'(x, t)|^{\alpha+2}}{a(u(x, t))} dx \leq 0,$$

where

$$A_1 = 2\bar{k}_2 - \frac{k_0^{\frac{(\alpha-\beta)}{(\alpha-1)}}(\alpha-1)}{(\alpha-\beta)} \left( \frac{(\alpha-\beta)}{(1-\beta)} \bar{k}_2 \right)^{-\frac{(1-\beta)}{(\alpha-1)}} > 0.$$

Inequality (3.8) shows that  $\Phi$  is a monotonically decreasing and bounded function for all  $t \geq 0$ . We define

$$\Theta(t) = [\Phi(t) - \Phi(t+1)]^{\frac{1}{\beta+2}}, \quad t \geq 0.$$

Integrating (3.8) we get

$$(3.9) \quad \Phi(t) + \frac{2c_1}{\rho} \int_0^t |\delta'(\tau)|^2 d\tau + A_1 \int_0^t \int_0^1 \frac{|u'(x, \tau)|^{\beta+2}}{a(u(x, \tau))} dx d\tau \\ + \bar{k}_2 \int_0^t \int_0^1 \frac{|u'(x, \tau)|^{\alpha+2}}{a(u(x, \tau))} dx d\tau \leq \Phi(0),$$

and

$$(3.10) \quad \int_t^{t+1} |\delta'(\tau)|^2 d\tau + \int_t^{t+1} \|u'(\tau)\|_{L^{\beta+2}(0,1)}^{\beta+2} d\tau \\ + \int_t^{t+1} \|u'(\tau)\|_{L^{\alpha+2}(0,1)}^{\alpha+2} d\tau \leq A_2 \Theta^{\beta+2}(t), \quad \text{for all } t \geq 0.$$

By the mean value theorem for integrals, there exist  $t_1 \in [t, t + \frac{1}{4}]$  and  $t_2 \in [t + \frac{3}{4}, t + 1]$  such that

$$(3.11) \quad \|u'(t_i)\| \leq 2A_2^{\frac{1}{\beta+2}} \Theta(t), \quad |\delta'(t_i)| \leq 2A_2^{\frac{1}{2}} \Theta^{\frac{\beta+2}{2}}(t), \quad \text{for } i = 1, 2.$$

Moreover, from (3.6), we obtain

$$(3.12) \quad \|u(t)\|_V \leq \Phi^{\frac{1}{2}}(t) \quad \text{and} \quad |\delta(t)| \leq \sqrt{\frac{\rho}{c_0}} \Phi^{\frac{1}{2}}(t).$$

Another simple consequence of the mean value theorem is that there exists a point  $t^* \in [t_1, t_2]$  such that

$$(3.13) \quad \Phi(t^*) \leq 2 \int_{t_1}^{t_2} \Phi(\tau) d\tau.$$

Integrating (3.7) from  $t^*$  to  $t$ , we find

$$(3.14) \quad \Phi(t) = \Phi(t^*) + 2 \int_t^{t^*} \int_0^1 \frac{g(u'(x, \tau)) u'(x, \tau)}{a(u(x, \tau))} dx d\tau \\ + \frac{2c_1}{\rho} \int_t^{t^*} |\delta'(\tau)|^2 d\tau + \int_t^{t^*} \int_0^1 \frac{a'(u(x, \tau)) |u'(x, \tau)|^3}{[a(u(x, \tau))]^2} dx d\tau.$$

To estimate the right hand side of (3.14), we observe that (2.1), (3.1) and (3.10) give

$$(3.15) \quad \begin{aligned} & 2 \int_t^{t^*} \int_0^1 \frac{g(u'(x, \tau)) u'(x, \tau)}{a(u(x, \tau))} dx d\tau \\ & \leq \frac{2\bar{k}_1}{a_0} \left( \int_t^{t^*} \|u'(\tau)\|_{L^{\beta+2}(0,1)}^{\beta+2} d\tau + \int_t^{t^*} \|u'(\tau)\|_{L^{\alpha+2}(0,1)}^{\alpha+2} d\tau \right) \leq \frac{4\bar{k}_1 A_2}{a_0} \Theta^{\beta+2}(t); \end{aligned}$$

$$(3.16) \quad \frac{2c_1}{\rho} \int_t^{t^*} |\delta'(\tau)|^2 d\tau \leq \frac{2c_1 A_2}{a_0} \Theta^{\beta+2}(t);$$

$$(3.17) \quad \begin{aligned} & \int_t^{t^*} \int_0^1 \frac{a'(u(x, \tau)) |u'(x, \tau)|^3}{[a(u(x, \tau))]^2} dx d\tau \leq \frac{k_0}{a_0} \int_t^{t^*} \int_0^1 |u'(x, \tau)|^3 dx d\tau \\ & \leq A_3 \left( \int_t^{t^*} \|u'(\tau)\|_{L^{\beta+2}(0,1)}^{\beta+2} d\tau + \int_t^{t^*} \|u'(\tau)\|_{L^{\alpha+2}(0,1)}^{\alpha+2} d\tau \right) \leq A_4 \Theta^{\beta+2}(t). \end{aligned}$$

Taking into account (3.13) – (3.17), we find

$$(3.18) \quad \Phi(t) \leq 2 \int_{t_1}^{t_2} \Phi(\tau) d\tau + A_5 \Theta^{\beta+2}(t).$$

We need an estimate for the integral  $\int_{t_1}^{t_2} \Phi(\tau) d\tau$ . Using (3.6) and (3.10), we have

$$(3.19) \quad \int_{t_1}^{t_2} \Phi(\tau) d\tau \leq A_6 (\Theta^2(t) + \Theta^{\beta+2}(t)) + \int_{t_1}^{t_2} \|u(\tau)\|_V^2 d\tau + \frac{c_0}{\rho} \int_{t_1}^{t_2} |\delta(\tau)|^2 d\tau.$$

Multiplying (1.1) by  $u(x, t)$  and integrating over  $(0, 1)$  we get

$$\begin{aligned} \|u(t)\|_V^2 &= -\frac{d}{dt} \int_0^1 \frac{u'(x, t) u(x, t)}{a(u(x, t))} dx + \int_0^1 \frac{|u'(x, t)|^2}{a(u(x, t))} dx \\ &\quad - \int_0^1 \frac{a'(u(x, t)) |u'(x, t)|^2 u(x, t)}{[a(u(x, t))]^2} dx + \delta'(t) u(1, t) \\ &\quad - \int_0^1 \frac{g(u'(x, t)) u(x, t)}{a(u(x, t))} dx. \end{aligned}$$

Integrating this from  $t_1$  to  $t_2$ , we obtain

$$(3.20) \quad \begin{aligned} & \int_{t_1}^{t_2} \|u(\tau)\|_V^2 d\tau = \int_0^1 \frac{u'(x, t_1) u(x, t_1)}{a(u(x, t_1))} dx - \int_0^1 \frac{u'(x, t_2) u(x, t_2)}{a(u(x, t_2))} dx \\ & \quad + \int_{t_1}^{t_2} \int_0^1 \frac{|u'(x, \tau)|^2}{a(u(x, \tau))} dx d\tau - \int_{t_1}^{t_2} \int_0^1 \frac{a'(u(x, \tau)) |u'(x, \tau)|^2 u(x, \tau)}{[a(u(x, \tau))]^2} dx d\tau \end{aligned}$$

$$+ \int_{t_1}^{t_2} \delta'(\tau) u(1, \tau) d\tau - \int_{t_1}^{t_2} \int_0^1 \frac{g(u'(x, \tau)) u(x, \tau)}{a(u(x, \tau))} dx d\tau.$$

Since  $V \hookrightarrow C([0, 1])$ , making use of (3.10) – (3.12), we can estimate each term of the right hand side of (3.20) for all  $\eta > 0$  as follows:

$$(3.21) \quad \int_0^1 \frac{u'(x, t_i) u(x, t_i)}{a(u(x, t_i))} dx \leq \frac{2A_2^{\frac{1}{\beta+2}}}{a_0} \Theta(t) \Phi^{\frac{1}{2}}(t) \\ \leq \eta \left( \sup_{t \leq \tau \leq t+1} \Phi(\tau) \right) + C(\eta) \Theta^2(t); \quad \text{for } i = 1, 2.$$

$$(3.22) \quad \int_{t_1}^{t_2} \int_0^1 \frac{|u'(x, \tau)|^2}{a(u(x, \tau))} dx d\tau \leq \frac{A_2^{\frac{2}{\beta+2}}}{a_0} \Theta^2(t);$$

$$(3.23) \quad \int_{t_1}^{t_2} \int_0^1 \frac{a'(u(x, \tau)) |u'(x, \tau)|^2 u(x, \tau)}{[a(u(x, \tau))]^2} dx d\tau \\ \leq \frac{k_0}{a_0} \int_{t_1}^{t_2} \left( \max_{0 \leq x \leq 1} |u(x, \tau)| \right) \int_0^1 |u'(x, \tau)|^2 dx d\tau \\ \leq A_7 \int_{t_1}^{t_2} \|u(\tau)\|_V \|u'(\tau)\|^2 d\tau \leq A_8 \Theta^2(t);$$

$$(3.24) \quad \int_{t_1}^{t_2} \delta'(\tau) u(1, \tau) d\tau \leq \int_{t_1}^{t_2} \left( \max_{0 \leq x \leq 1} |u(x, \tau)| \right) |\delta'(\tau)| d\tau \\ \leq A_9 \int_{t_1}^{t_2} \|u(\tau)\|_V |\delta'(\tau)| d\tau \leq \eta \left( \sup_{t \leq \tau \leq t+1} \Phi(\tau) \right) + C(\eta) \Theta^{\beta+2}(t);$$

$$(3.25) \quad \int_{t_1}^{t_2} \int_0^1 \frac{g(u'(x, \tau)) u(x, \tau)}{a(u(x, \tau))} dx d\tau \\ \leq \frac{\bar{k}_1}{a_0} \int_{t_1}^{t_2} \int_0^1 \left( |u'(x, \tau)|^{\beta+1} + |u'(x, \tau)|^{\alpha+1} \right) |u(x, \tau)| dx d\tau \\ \leq \eta \int_{t_1}^{t_2} \left[ \max_{0 \leq x \leq 1} |u(x, \tau)|^\beta + \max_{0 \leq x \leq 1} |u(x, \tau)|^\alpha \int_0^1 |u(x, \tau)|^2 dx \right] d\tau \\ + C(\eta) \int_{t_1}^{t_2} \left( \|u'(\tau)\|_{L^{\beta+2}(0,1)}^{\beta+2} + \|u'(\tau)\|_{L^{\alpha+2}(0,1)}^{\alpha+2} \right) d\tau \\ \leq C \eta \left( \sup_{t \leq \tau \leq t+1} \Phi(\tau) \right) + C(\eta) \Theta^{\beta+2}(t).$$

From (3.20) – (3.25), we have

$$(3.26) \quad \int_{t_1}^{t_2} \|u(\tau)\|_V^2 ds \leq 4\eta \left( \sup_{t \leq \tau \leq t+1} \Phi(\tau) \right) + A(\eta) \left( \Theta^2(t) + \Theta^{\beta+2}(t) \right).$$

It remains to estimate the last term of (3.19). We multiply (1.3) by  $\delta(t)$  and integrate from  $t_1$  to  $t_2$ , to obtain

$$(3.27) \quad \frac{c_0}{\rho} \int_{t_1}^{t_2} |\delta(\tau)|^2 d\tau = u(1, t_1)\delta(t_1) - u(1, t_2)\delta(t_2) + \int_{t_1}^{t_2} u(1, \tau)\delta'(\tau) d\tau \\ + \frac{1}{\rho}\delta'(t_1)\delta(t_1) - \frac{1}{\rho}\delta'(t_2)\delta(t_2) + \frac{1}{\rho} \int_{t_1}^{t_2} |\delta'(\tau)|^2 d\tau - \frac{c_1}{\rho} \int_{t_1}^{t_2} \delta(\tau)\delta'(\tau) d\tau.$$

Let us estimate the right hand side of (3.27). First of all, using (3.12), we can see that

$$(3.28) \quad |u(1, t_i)\delta(t_i)| \leq \bar{c} \sqrt{\frac{\rho}{c_0}} \Phi(t), \quad \text{for } i = 1, 2.$$

Here  $\bar{c}$  is given by (3.4), and we have used that  $V \hookrightarrow C([0, 1])$  and the fact that  $\Phi$  is a monotonically decreasing function.

On the other hand,

$$(3.29) \quad \frac{1}{\rho}\delta'(t_i)\delta(t_i) \leq \eta \left( \sup_{t \leq \tau \leq t+1} \Phi(\tau) \right) + C(\eta) \Theta^{\beta+2}(t), \quad \text{for } \eta > 0, \quad i = 1, 2;$$

$$(3.30) \quad \frac{1}{\rho} \int_{t_1}^{t_2} |\delta'(\tau)|^2 d\tau \leq \frac{A_2}{\rho} \Theta^{\beta+2}(t);$$

$$(3.31) \quad \frac{c_1}{\rho} \int_{t_1}^{t_2} \delta(\tau)\delta'(\tau) d\tau \leq \frac{\epsilon c_0}{\rho} \int_{t_1}^{t_2} |\delta(\tau)|^2 d\tau + C(\epsilon) \Theta^{\beta+2}(t), \quad \text{for all } \epsilon \geq 0.$$

Taking into account (3.24), (3.27) and (3.28) – (3.31), we get

$$\frac{c_0}{\rho}(1 - \epsilon) \int_{t_1}^{t_2} |\delta(\tau)|^2 d\tau \leq 2\bar{c} \sqrt{\frac{\rho}{c_0}} \Phi(t) + 3\eta \left( \sup_{t \leq \tau \leq t+1} \Phi(\tau) \right) \\ + (3C(\eta) + \frac{A_2}{\rho} + C(\epsilon)) \Theta^{\beta+2}(t),$$

choosing  $\epsilon = \frac{1}{2}$ , we find

$$(3.32) \quad \frac{c_0}{\rho} \int_{t_1}^{t_2} |\delta(\tau)|^2 d\tau \leq 4\bar{c} \sqrt{\frac{\rho}{c_0}} \Phi(t)$$

$$+6\eta \left( \sup_{t \leq \tau \leq t+1} \Phi(\tau) \right) + (6C(\eta) + A_{10}) \Theta^{\beta+2}(t).$$

From (3.18), (3.19), (3.26) and (3.32), it follows

$$(1 - 8\bar{c} \sqrt{\frac{\rho}{c_0}}) \Phi(t) \leq A_{11}\eta \left( \sup_{t \leq \tau \leq t+1} \Phi(\tau) \right) + \bar{A}(\eta) (\Theta^2(t) + \Theta^{\beta+2}(t)).$$

This inequality and assumption (3.4), for  $\eta > 0$  sufficiently small, imply

$$\sup_{t \leq \tau \leq t+1} \Phi(\tau) \leq A_{12}(\Phi(0)) \Theta^2(t), \quad \text{for all } t \geq 0,$$

which can be rewritten as

$$(3.33) \quad \sup_{t \leq \tau \leq t+1} \Phi^{1+\frac{\beta}{2}}(\tau) \leq C(\Phi(0)) [\Phi(t) - \Phi(t+1)] \quad \text{for all } t \geq 0.$$

By lemma of M. Nakao ( see [6] and [7] ), (3.5) holds, there with theorem 3.1 is proved.

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